

# High-Power 50-GHz Double-Drift-Region IMPATT Oscillators with Improved Bias Circuits for Eliminating Low-Frequency Instabilities

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**Abstract**—Low-frequency instabilities in millimeter-wave double-drift-region (DDR) IMPATT diodes are investigated and new oscillator circuits with the improved bias circuits for eliminating the low-frequency instability are developed. DDR IMPATT diodes mounted in these circuits exhibited a maximum free-running oscillation power of 1.6 W at 55.5 GHz with 11.5-percent conversion efficiency. A highly stabilized oscillator was also constructed with the maximum output power of 1 W and the frequency stability 0.3 ppm/mA at 51.86 GHz.

## INTRODUCTION

THE double-drift-region (DDR) IMPATT diodes are presently considered as the only solid-state active device that can produce output power on the order of 1 W in the millimeter-wave range. It has been confirmed that the millimeter-wave output power and efficiency are significantly higher for DDR diodes than that obtained with the single-drift-region (SDR) IMPATT diodes [1]. DDR diodes, however, have not yet been practically utilized at their highest power and efficiency levels. The reason for this is that there still remain some problems to be solved in the practical utilization of DDR diodes. One of these problems is that DDR diodes are prone to low-frequency instabilities because of the high conversion efficiency [2].

Low-frequency instabilities in IMPATT diodes have a detrimental effect on microwave oscillation [3]. Low-frequency instabilities would result in excessive noise increase, tuning-induced diode burnout, and microwave output power decrease. These instabilities are more prominent in higher efficiency, higher power, and higher frequency diodes. That is, these phenomena are serious problems in millimeter-wave IMPATT diodes, especially in millimeter-wave DDR IMPATT diodes which have a high power capability. Therefore, the elimination of the low-frequency instability is essential in the practical utilization of millimeter-wave DDR diodes. One can expect that high-power millimeter-wave oscillators having low sideband noise can be obtained if this low-frequency instability is suppressed.

In this paper, the low-frequency instabilities in millimeter-wave DDR IMPATT diodes are investigated and the low-frequency impedance of DDR diodes is derived. It is shown that DDR diodes are more prone to instability than

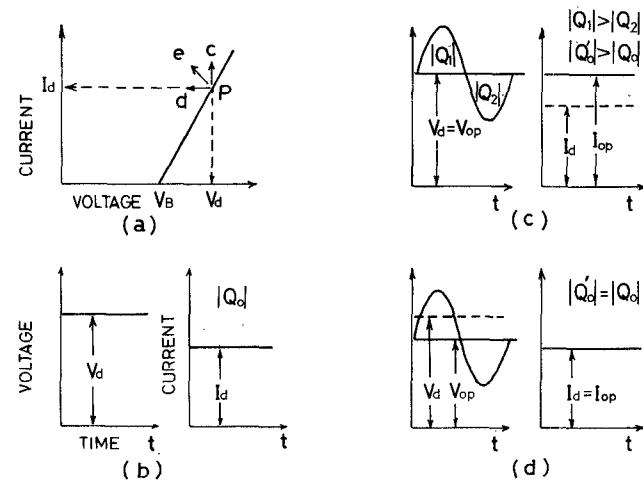


Fig. 1. Origins of the low-frequency negative resistance in the IMPATT diode. (a) Voltage-current characteristics of the reverse-biased IMPATT diode. (b) RF voltage amplitude  $V_{RF} = 0$ .  $|Q_0|$  is the amount of the charge induced by avalanche. (c)  $V_{RF} > 0$ . The amount of charge  $|Q_1|$  induced between 0 and  $\pi$  in the phase of RF voltage becomes larger than that of the charge  $|Q_2|$  which is removed between  $\pi$  and  $2\pi$  because of the excess increase of the charge induced by avalanche due to the extreme nonlinearity of the ionization coefficient. (d)  $V_{RF} > 0$ . When the constant current is applied to the diode, the operating voltage  $V_{op}$  decreases so as to satisfy the condition of  $|Q_0'| = |Q_0|$ , where  $|Q_0'|$  is the amount of the charge induced by avalanche in case of  $V_{RF} > 0$ .

SDR diodes. Some new oscillator circuits with improved bias circuits for eliminating low-frequency instabilities are proposed. DDR IMPATT diodes mounted in these circuits exhibited output powers of 1.6 W at 55.5 GHz with 11.5-percent conversion efficiency. A highly stabilized oscillator was also constructed with the maximum output power of 1 W at 51.86 GHz and the frequency stability 0.3 ppm/mA.

## ORIGINS OF LOW-FREQUENCY NEGATIVE RESISTANCE IN IMPATT DIODES AND INFLUENCES OF THE LOW-FREQUENCY INSTABILITY

Low-frequency negative resistance originates in the excessive increase of the charges induced by avalanche with increasing RF voltage amplitude  $V_{RF}$  due to the extreme nonlinearity of the ionization coefficient  $\alpha$  as a function of the electric field  $E$  [4]. This mechanism is explained in more detail in Fig. 1. Fig. 1(a) shows the voltage-current characteristics of the reverse-biased IMPATT diode. At point P, in the case of  $V_{RF} = 0$ , the applied voltage  $V$  and

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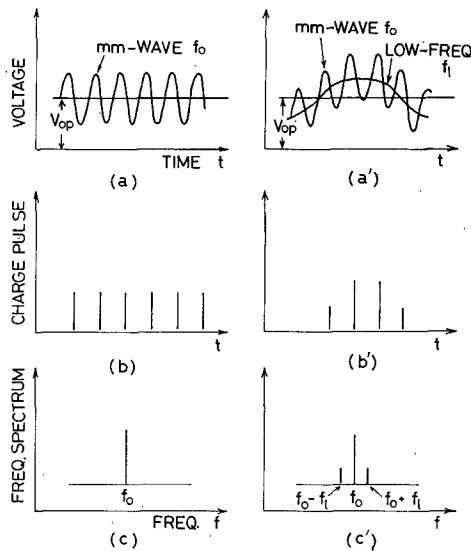


Fig. 2. The outline of the influence of the low-frequency instability on the IMPATT operation of the diode. The left-hand side represents the normal IMPATT oscillation and the right-hand side the anomalous oscillation where low-frequency instability appears.

the direct current  $I_d$  is constant with time, as shown in Fig. 1(b), where  $|Q_0|$  is the amount of the charge induced by avalanche. With the RF voltage applied [in Fig. 1(c)], the amount of the charge  $|Q_1|$  induced between 0 and  $\pi$  in the phase of the RF voltage becomes larger than that of the charge  $|Q_2|$  removed between  $\pi$  and  $2\pi$  due to the nonlinearity of  $\alpha$ . That is,  $|Q_1| > |Q_2|$  so that the amount of the charge  $|Q_0'|$  induced in one cycle of the RF voltage is larger than  $|Q_0|$ . Then, if the applied voltage  $V_d$  is constant, the operating current  $I_{op}$  increases more than  $I_d$  [as shown in Fig. 1(a), arrow c]. If the constant current is applied to the diode [in Fig. 1(d)], the operating voltage  $V_{op}$  decreases so as to satisfy the condition of  $|Q_0'| = |Q_0|$  [in Fig. 1(a), arrow d]. If both  $V_{op}$  and  $I_{op}$  are not constrained,  $V_{op}$  decreases and  $I_{op}$  increases with increasing  $V_{RF}$ , i.e.,  $dV_{op}/dI_{op}$  is negative. This is the reason why the low-frequency negative resistance appears in IMPATT diodes.

The outline of the influence of low-frequency instability on the IMPATT operation is shown in Fig. 2. The left-hand side of Fig. 2 represents the normal IMPATT operation and the right-hand side the anomalous one where the low-frequency instability appears. When the low-frequency voltage amplitude superimposes on the millimeter-wave one, the phase and the amount of the charge pulse generated by avalanche will shift randomly, whereas a constant charge pulse train is generated in the  $\pi$  phase of the millimeter-wave voltage amplitude when the diode operates in the normal IMPATT mode. It is assumed from these speculations that the millimeter-wave output power will be reduced by the occurrence of the low-frequency oscillation.

#### LOW-FREQUENCY IMPEDANCE OF DDR IMPATT DIODES AND ELIMINATION OF LOW-FREQUENCY INSTABILITIES

The equation representing the relation among the operating voltage  $V_{op}$ , the applied current  $I_d$ , and the RF voltage amplitude  $V_{RF}$  in SDR IMPATT diodes was first derived

by Read [5]

$$V_{op} = R_{sc} I_d - \frac{m}{4WE_c} V_{RF}^2 + \text{constant} \quad (1)$$

where  $R_{sc}$  is the space charge resistance,  $W$  is the depletion layer width,  $E_c$  is the critical field for avalanche breakdown,  $m$  is  $(E_c/\alpha) \cdot (d\alpha/dE)|_{E_c}$ , and  $\alpha$  is the ionization coefficient.

A DDR diode has equal drift layers for holes and electrons. As these layers are arranged in series, the operating voltage  $V_{op}$  in the DDR diode is composed of the algebraic sum of the operating voltage  $V_{op(p)}$  applied to p layer and  $V_{op(n)}$  to n layer, i.e.,  $V_{op} = V_{op(p)} + V_{op(n)}$ . Equations similar to (1) can be derived for p and n layers in DDR diodes. By adding two equations, we can have the following equation for DDR diodes:

$$V_{op} = R_{sc} I_d - \frac{m_n + m_p}{8WE_c} V_{RF}^2 + \text{constant} \quad (2)$$

where  $m_n$  and  $m_p$  is  $(E_c/\alpha_n) \cdot (d\alpha_n/dE)|_{E_c}$  and  $(E_c/\alpha_p) \cdot (d\alpha_p/dE)|_{E_c}$ , respectively,  $\alpha_n$  and  $\alpha_p$  are ionization coefficients for electron and hole, respectively, and  $W$  is the total active layer width in DDR diodes. In addition, the following assumptions have been made:

$$W_{(p)} = W_{(n)} = 1/2 \cdot W \quad V_{op(p)} = V_{op(n)} = 1/2 \cdot V_{op}$$

$$V_{RF(p)} = V_{RF(n)} = 1/2 \cdot V_{RF} \quad R_{sc(p)} + R_{sc(n)} = R_{sc}$$

where the suffixes (p) and (n) represent the p- and n-type layers, respectively.

Differentiating both sides of (2) with respect to  $I_d$ , we obtain

$$\frac{dV_{op}}{dI_d} = R_{sc} - \frac{m_n + m_p}{4WE_c} \cdot V_{RF} \cdot \left( \frac{dV_{RF}}{dI_d} \right). \quad (3)$$

Using the relation  $P_{out} = 1/2 \cdot |G_{op}| \cdot V_{RF}^2$ , (3) can be written

$$\frac{dV_{op}}{dI_d} = R_{sc} - \frac{m_n + m_p}{4WE_c} \cdot \frac{1}{|G_{op}|} \cdot \left( \frac{dP_{out}}{dI_d} \right) \quad (4)$$

where  $G_{op}$  is the operation conductance and  $P_{out}$  is the output power. If we define  $dV_{op}/dI_d \equiv R_{DDR}$ ,  $R_{DDR}$  represents the low-frequency negative resistance of DDR diodes.

Brackett [3] derived the low-frequency impedance  $Z_{SDR}$  for SDR diodes given by

$$Z_{SDR} = R_{sc} - \frac{R_-}{1 + j \frac{\omega}{\gamma}} \quad (5a)$$

$$R_- = \frac{m}{2WE_c} \cdot \frac{1}{|G_{op}|} \cdot \left( \frac{dP_{out}}{dI_d} \right) \quad (5b)$$

where  $\gamma$  is a quantity associated with the saturation parameters of the large-signal behavior of the diode, the millimeter-wave oscillation frequency, and the loaded cavity  $Q$  [3]. The second term on the right-hand side of (4) is similar to (5b). Hence the low-frequency impedance  $Z_{DDR}$

TABLE I  
PHYSICAL DATA AND NEGATIVE RESISTANCE VALUES FOR MILLIMETER-WAVE Si-SDR AND -DDR IMPATT'S

	Si SDR	Si DDR	Source
Diode	37-1H	39-16S	
Diode Capacitance $C_0$ (pF)	1.4	1.2	Measured
C breakdown (pF)	0.32	0.28	Measured
Junction Area $S$ (cm <sup>2</sup> )	$1.6 \times 10^{-5}$	$2.2 \times 10^{-5}$	Calculated
Series Resistance $R_s$ (ohms)	1.06	0.85	Estimated <sup>(6)</sup>
Breakdown Voltage $V_B$ (volts)	17.3	21.3	Measured
Depletion Layer Width $W$ (microns)	0.72 at 23(V)	0.89 at 29(V)	Estimated
Critical Field $E_c$ (volts/cm)	$5.5 \times 10^{-5}$	$4.7 \times 10^{-5}$	Estimated
$\frac{E_c}{\alpha} \left( \frac{d\alpha}{dE} \right) E_c$	$\alpha_n = 1.38 \times 10^5$ (cm <sup>-1</sup> ) $\alpha_n = 9.20 \times 10^{-1}$ (V <sup>-1</sup> ) $m_n = 3.67$	$\alpha_n = 9.20 \times 10^4$ (cm <sup>-1</sup> ) $\alpha_p = 2.22 \times 10^4$ (cm <sup>-1</sup> ) $\alpha_p = 9.20 \times 10^{-1}$ (V <sup>-1</sup> ) $\alpha_p = 3.48 \times 10^{-1}$ (V <sup>-1</sup> ) $m_n = 4.72$ $m_p = 7.40$	Calculated <sup>(4)</sup>
$\frac{dP_{out}}{dI_d}$ (watts/amp)	4.0	5.9	Measured
Negative Conductance $G_{op}$ (mhos)	$10.7 \times 10$ at 457 mW	$S = 1.6 \times 10^{-5}$ cm <sup>2</sup> $7.4 \times 10^{-3}$	$S = 2.2 \times 10^{-5}$ cm <sup>2</sup> $10.2 \times 10^{-3}$ at 1023 mW
$R_-$ (ohms)	17.3	57.5	41.7
$R_{sc}$ (ohms)	5.0	8.0	6.0
Low Frequency Negative Resistance at $\omega = 0$ (ohms)	-12.3	-49.5	-35.7
Loaded Cavity Q $Q_L$	30	38	38
Millimeter-wave Oscillation Frequency $f_0$ (GHz)	55.5	55.5	55.5
Maximum Frequency $f_{max}$ (GHz)	3.4	3.6	3.6
			0.34

of DDR diodes can be expressed by

$$Z_{DDR} = R_{sc} - \frac{R_-}{1 + j \frac{\omega}{\gamma}} \quad (6a)$$

$$R_- = \frac{m_n + m_p}{4WE_c} \cdot \frac{1}{|G_{op}|} \cdot \left( \frac{dP_{out}}{dI_d} \right). \quad (6b)$$

At  $\omega = 0$ ,  $Z_{DDR} = R_{DDR}$  and using each parameter of DDR (39-16S) as listed in Table I,  $R_{DDR}$  is calculated to be about (-) 50 Ω, which is four times as large as  $R_{SDR}$ . This shows that these instabilities will be observed more frequently in DDR than in SDR.

We can derive from (6a) the maximum frequency  $f_m$  for which  $Z_{DDR}$  has a negative real part. The maximum frequency  $f_{max}$  is given by (see [3])

$$f_{max} = \frac{f_0}{Q_L} \sqrt{\frac{R_-}{R_{sc}} - 1} \quad (7)$$

where  $f_0$  is the millimeter-wave oscillation frequency, and  $Q_L$  is the loaded cavity  $Q$ . The calculated value of  $f_{max(DDR)}$  for DDR diodes (39-16S) using (7) is about 3.6 GHz as listed in Table I, while  $f_{max(SDR)}$  for SDR diodes is about 3.4 GHz. Smith chart plots of (-)  $Z_{DDR}$  and (-)  $Z_{SDR}$  with the frequency  $\omega$  are shown in Fig. 3(a) in the case where the junction area is  $1.6 \times 10^{-5}$  cm<sup>2</sup> for both SDR and DDR.

The  $f_{max}$  in millimeter-wave DDR diodes is several hundred times as large as that in microwave diodes. For example, the  $f_{max}$  is about 16 MHz in J-band GaAs IMPATT diodes [3]. The suppression of instabilities in higher frequency diodes is more difficult than in lower frequency

diodes. Therefore, these phenomena are more prominent in higher frequency diodes.

The outline of the frequency dependence of the negative resistance which DDR (39-16S) can have is illustrated in Fig. 3(b). Region 1 shows the intrinsic negative resistance area of DDR diodes, where  $\omega_a$  is the avalanche frequency. Region 2 shows the frequency band in which the RF voltage-induced negative resistance appears. The dotted line means that thermal effects become important and introduce further positive resistance. The low-frequency negative resistance (region 2) appears only when the millimeter-wave oscillation builds up (region 1). The load for this low-frequency negative resistance is the bias-circuit impedance seen from the diode. Therefore, in order to suppress the low-frequency instability in DDR diodes, we consider the following two methods (conditions):

1) A method for making the bias-circuit impedance extremely small (~ short circuit) compared with (-)  $Z_{DDR}$  at frequencies in excess of  $f_{min(DDR)}$  [see Fig. 3(b)] for the low-frequency instabilities.

2) A method for making the bias-circuit impedance greater than (-)  $Z_{DDR}$  at frequencies lower than  $f_{max(DDR)}$  for DDR. (This is the absolutely stable condition for low-frequency oscillation.)

In our experiments the low-frequency instability was observed to be prominent in the frequency range from 100 MHz<sup>1</sup> to 1 GHz. Therefore, if the bias-circuit impedance is kept small in this frequency range, condition 1) will be

<sup>1</sup> The thermal time constant should be of the order of at least 1 μs. However, the  $f_{min}$  in our experiments was about 100 MHz of which time constant was 1/100 μs. The reason why this discrepancy occurs is still unclear.

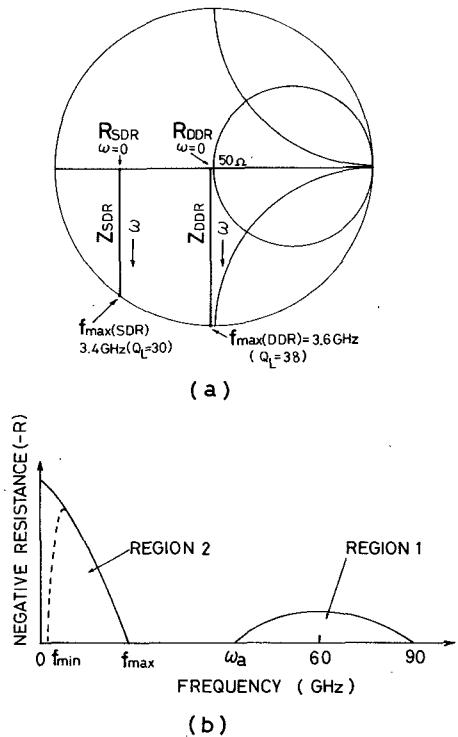


Fig. 3. (a) Smith chart plots of the low-frequency impedance  $(-Z_{DDR}(\omega))$  and  $(-Z_{SDR}(\omega))$  for DDR and SDR IMPATT diodes, respectively. Equation (6) gives a slight curve on the Smith chart when normalized by  $Z_0 = 50 \Omega$  and gives a straight vertical line when normalized by  $R_{sc}$  (which is usually smaller than  $50 \Omega$ ). In this case, for simplicity, (6) has been drawn as the straight vertical line even if normalized by  $Z_0 = 50 \Omega$ . It does not influence the result of this discussion. The junction area is  $1.6 \times 10^{-5} \text{ cm}^2$  for both SDR and DDR. (b) Negative resistance of DDR diodes versus the frequency. Region 1 shows the intrinsic negative resistance area of DDR diodes, where  $\omega_a$  is the avalanche frequency. Region 2 shows the frequency band of the RF voltage-induced negative resistance. The dotted line means that in a very-low-frequency band, thermal effects become important and introduce further positive resistance. The low-frequency negative resistance appears only when the millimeter-wave frequency oscillation builds up.

satisfied. It is seen from Fig. 3(a) that suppressing the instability in DDR diodes is more difficult than in SDR diodes, since  $Z_{DDR}$  and  $f_{max(DDR)}$  are greater than  $Z_{SDR}$  and  $f_{max(SDR)}$ , respectively.

#### BIAS-CIRCUIT IMPEDANCE OF THE CONVENTIONAL COAXIAL-WAVEGUIDE CIRCUIT AND CAP-TYPE CIRCUIT

The bias-circuit impedance  $Z_B(\omega)$  seen from the diode was measured with a network analyzer and a small coaxial probe with  $Z_0 = 50 \Omega$  [9], in place of the diode. Fig. 4(b) shows the bias-circuit impedance locus  $Z_{B1}(\omega)$  of the conventional coaxial-waveguide circuit as illustrated in Fig. 4(a), which has been used in millimeter-wave communication systems [10]. It is seen from Fig. 4(b) that the  $Z_{B1}(\omega)$  locus intersects the  $(-Z_{DDR}(\omega))$  locus within frequencies from several tens of megahertz to about 1 GHz. In this frequency band the low-frequency instabilities have been experimentally observed to be more prominent. The junction area of the diode used in this experiment is  $2.2 \times 10^{-5} \text{ cm}^2$  (see Table I). Fig. 5 shows the impedance locus of the same circuit as that of Fig. 4(a) except for the insertion of a parallel circuit of a resistor ( $R \sim 500 \Omega$ ) and an inductor ( $L \sim 23 \mu\text{H}$ ).

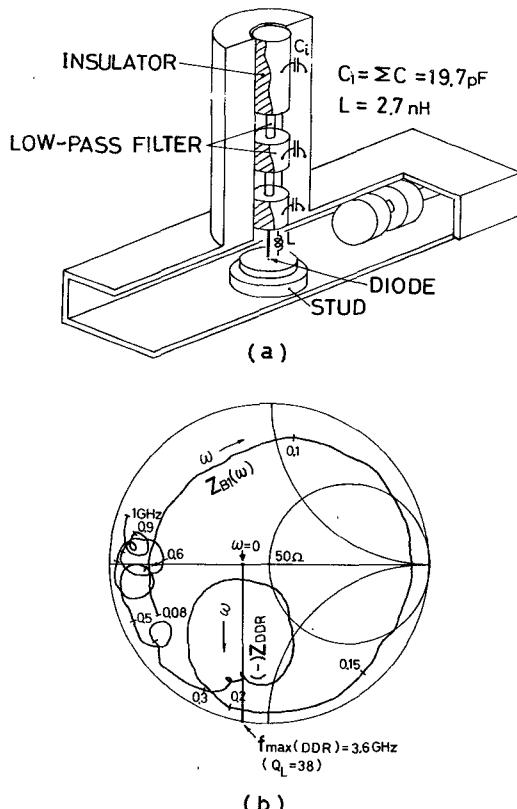


Fig. 4. (a) Coaxial-waveguide circuit used in millimeter-wave communication systems [10]. (b) The bias-circuit impedance locus of this circuit. The  $Z_{B1}(\omega)$  locus intersects the  $(-Z_{DDR}(\omega))$  locus. The junction area of the diode used in this experiment is  $2.2 \times 10^{-5} \text{ cm}^2$ .

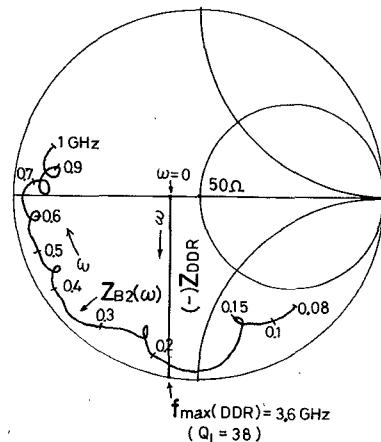


Fig. 5. The bias-circuit impedance locus  $Z_{B2}(\omega)$  of the same circuit as that of Fig. 4(a), except for the insertion of a parallel circuit of a resistor ( $R \sim 500 \Omega$ ) and an inductor ( $L \sim 23 \mu\text{H}$ ).

an inductor ( $L \sim 23 \mu\text{H}$ ) in series with the dc bias feed line between the diode cavity and the dc power supply. The  $Z_{B2}(\omega)$  locus is simpler at frequencies below 0.3 GHz than that of Fig. 4(b). But as this locus also intersects the  $(-Z_{DDR}(\omega))$  locus, no improvement can be expected for the low-frequency instability. At 0.65 GHz, the  $Z_{B2}(\omega)$  locus intersects the line of  $jX = 0$ . The reason for this is that the capacitance  $C_1$  [ $\sim 19.7 \text{ pF}$ , see Fig. 4(a)] and inductance  $L$  of the inner conductor in the coaxial line are in series

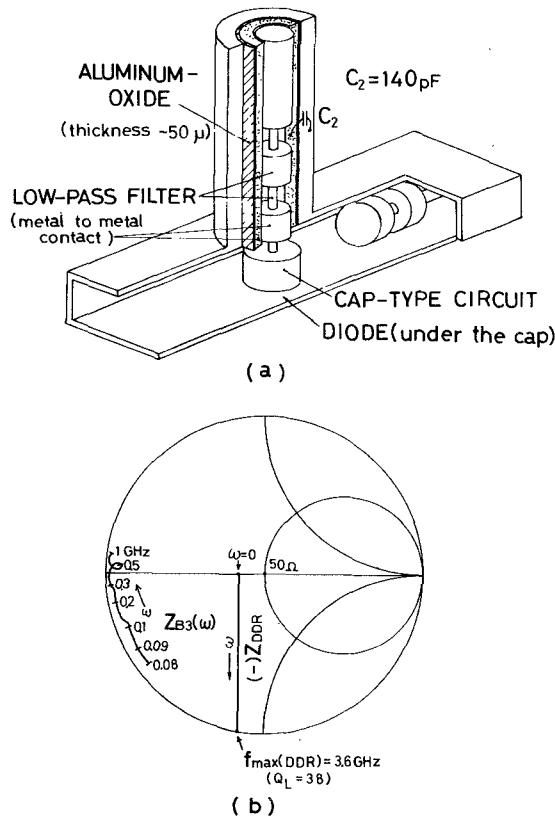


Fig. 6. (a) Improved cap-type circuit. A large capacitance ( $\sim 140$  pF) made of aluminum oxide (thickness  $\sim 50$   $\mu$ m) is placed so as to surround the low-pass filter. (b) The bias-circuit impedance locus  $Z_{B3}(\omega)$  of this circuit. Compared with Fig. 5, the impedance is successfully kept low at frequencies in excess of 80 MHz.

resonance at 0.65 GHz. The inductance  $L$  calculated from  $L = 1/\omega^2 C$  is about 2.7 nH, which is in reasonable agreement with the calculated value of 1.8 nH. It may be concluded that due to the capacitance of the low-pass filter, the  $Z_{B2}(\omega)$  locus represents the negative reactance at lower frequencies than 0.65 GHz.

As the millimeter-wave circuit satisfying condition 2) can not be easily constructed because of the very small dimensions, we first tried to realize method 1). In order to realize method 1), it is necessary to remove inductance  $L$  and to put a large capacitance in the vicinity of the diode. Therefore, a millimeter-wave circuit having an inductance near the diode, such as a coaxial-waveguide type, is not suitable for DDR IMPATT diodes. From this point of view, the cap-type circuit might be preferable for DDR diodes as it has no inductance near the diode. The bias-circuit impedance locus of the unimproved cap-type circuit, with the parallel circuit in the dc bias feed line at frequencies lower than 0.65 GHz, was observed to be similar to that of Fig. 5. Under this condition the low-frequency instability would not be suppressed. It is necessary to improve the cap-type circuit so as to satisfy condition 1). In order to make the bias-circuit impedance extremely small without producing any loss in millimeter-wave power and efficiency, a large capacitance ( $C_2 \sim 140$  pF) made of aluminum oxide (thickness 50  $\mu$ m) was placed so as to surround the low-pass filter as illustrated in Fig. 6(a). The bias-circuit impedance

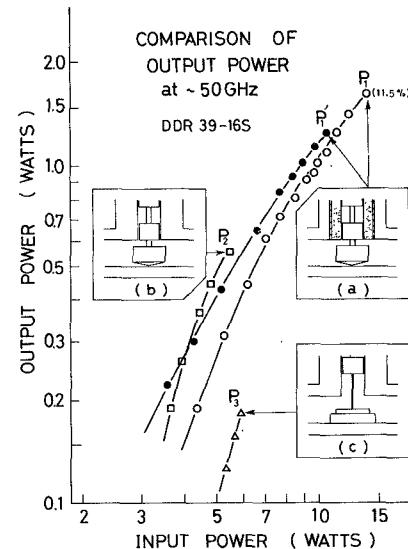


Fig. 7. Output power and conversion efficiency versus input power for typical devices mounted in the three kinds of circuits. The junction area of the diodes used is about  $2.2 \times 10^{-5}$   $\text{cm}^2$  and the typical value of the thermal resistance of these diodes is about  $25^\circ\text{C/W}$ . (a)  $\circ$ : improved cap-type circuit. Maximum power and efficiency obtained is  $\sim 1.6$  W and 11.5 percent, respectively. (b)  $\square$ : unimproved cap-type circuit. The low frequency instability occurred at  $P_2$ . (c)  $\triangle$ : conventional coaxial-waveguide circuit. The low-frequency instability occurred at  $P_3$ .

locus  $Z_{B3}(\omega)$  of this circuit is shown in Fig. 6(b). Compared with Fig. 5, the impedance was successfully kept low at frequencies in excess of 80 MHz. It is expected that the low-frequency instability will not build up in this circuit configuration.

#### EXPERIMENTAL RESULTS AND PRACTICAL DESIGN CONSIDERATIONS

The fabrication technique of diodes used in microwave measurements is the following: The diode pellet was mounted on a diamond heat sink and the junction area was controlled to be about  $2.2 \times 10^{-5}$   $\text{cm}^2$  by chemical etching. The typical value of the thermal resistance of these diodes is  $25^\circ\text{C/W}$ . The contact resistance of these diodes is fully reduced by forming the  $n^{++}$  ( $\sim 10^{20}$   $\text{cm}^{-3}$ ) layer near the surface of the  $n^+$  substrate by arsenic ion implantation [6].

Fig. 7 shows the output power and conversion efficiency versus the input power for typical devices mounted in the three kinds of circuits. In the case of the conventional coaxial-waveguide circuit as illustrated in insert (c), the oscillation frequency spectrum observed with the spectrum analyzer showed a large amount of sideband noise and the output power obtained was at most 200 mW. Next, using the unimproved cap-type circuit as illustrated in insert (b), the maximum output power obtained was 500 mW. When the input power was further increased, the low-frequency instability occurred and the output power did not increase any more.

In the case of the improved cap-type circuit as illustrated in insert (a), power and efficiency as high as 1.6 W (point  $P_1$ ) and 11.5 percent, respectively, were achieved at 55.5 GHz, which is the best power times frequency squared product ( $\text{pf}^2$ ) reported so far for any microwave solid-state

devices. At this power level, no degradation in frequency spectrum was observed. When the bias current was further increased, the diode burned out.

It is evident that DDR IMPATT diodes can exhibit their intrinsic high-power characteristics if the low-frequency instabilities are suppressed. In order to use these diodes in a practical system, it is necessary that the low-frequency instabilities be completely eliminated under every operating condition. This requires that the bias circuit be designed by method 2) to satisfy the absolutely stable condition. In addition, the circuit should be easily designed to operate the DDR IMPATT diodes at any specified frequency with the maximum power. The improved cap-type circuit described previously is not suitable for this purpose, because a slight change of bias current or load impedance will cause low-frequency instabilities at a high output power level. Furthermore, the design of the cap-type circuit for a specific frequency is not easy. Therefore, it is necessary to develop a new millimeter-wave oscillator circuit suitable for the practical application of DDR IMPATT diodes.

In order to construct a circuit which satisfies the absolutely stable condition, the bias-circuit impedance must be greater than  $(-)$   $Z_{DDR}$  in the frequencies below  $f_{max(DDR)}$ . For this purpose, either or both of the following conditions should be realized:

- 1) To make the circuit  $Q$  ( $Q_L$ ) so large that  $f_{max(DDR)}$  becomes lower than the frequency at which the  $Z_{B4}(\omega)$  locus intersects the  $(-)$   $Z_{DDR}$  locus.
- 2) To remove the capacitance in the vicinity of the diode and to insert an inductance or a parallel circuit of  $L$  and  $R$  in some place of the millimeter-wave circuit without affecting the millimeter-wave oscillation.

Brackett [3] has proposed that the low-frequency instability of a 6-GHz GaAs IMPATT could be eliminated by inserting a resistive choke stabilizing network in a bias feed coaxial post immediately above the diode. In the millimeter-wave range, however, because of the very small dimensions of the circuit, this method cannot be used successfully and results in the large degradation of the millimeter-wave power and efficiency.

#### CAVITY-STABILIZED DDR OSCILLATOR

From the aforementioned design considerations, the oscillator circuit shown in Fig. 8(a) was developed for 50-GHz-band applications using a WR-19 waveguide. A coaxial termination is incorporated in the bias circuit to give a high impedance loading to the diode in the frequency range of instabilities. To minimize the millimeter-wave dissipation, a radial line cavity band rejection filter is placed in the bias circuit. A high  $Q$  cylindrical cavity is placed in the waveguide circuit to increase the circuit  $Q$  of the oscillator. Millimeter-wave absorbers in the bias circuit and in the waveguide circuit are employed to suppress the spurious millimeter-wave oscillations. Impedance matching between the diode and the waveguide is made by a quarter-wave coaxial transformer which contacts the standoff of the diode. The oscillation frequency is determined by the geometry of this transformer and precise tuning can be made by the cylindrical cavity.

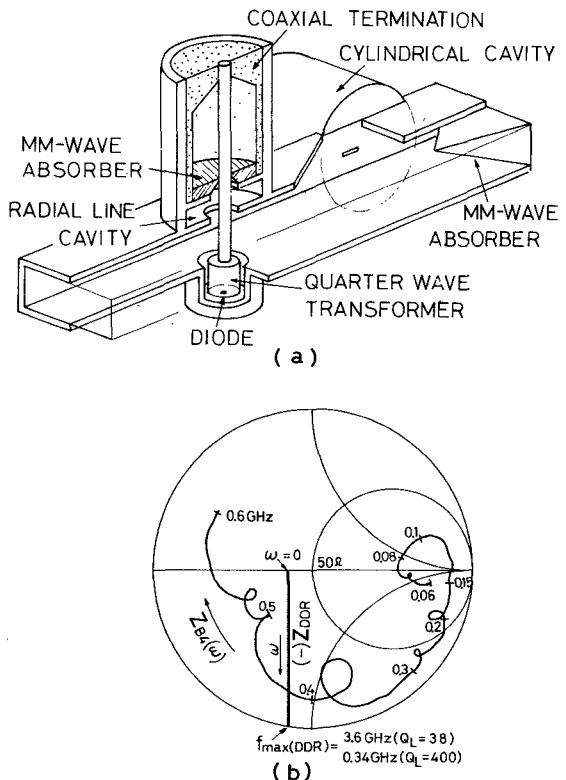


Fig. 8. (a) Cavity-stabilized DDR oscillator. The bias circuit provides a high-impedance coaxial termination to damp the instability. (b) The bias-circuit impedance locus of this circuit. The  $Z_{B4}$  locus intersects the  $(-)$   $Z_{DDR}(\omega)$  locus at 0.42 GHz.

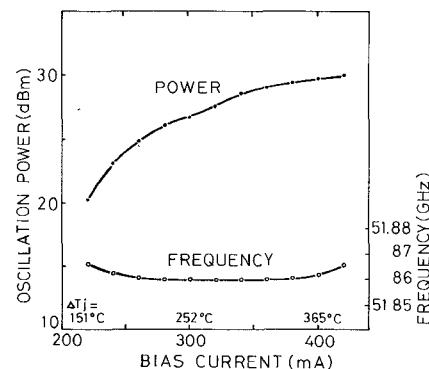


Fig. 9. Output power and oscillation frequency versus bias current for the highly stabilized DDR oscillator shown in Fig. 8(a). Maximum output power obtained is 1 W with the frequency stability less than 3 MHz over a 200-mA bias current range at 51.86 GHz.

The measured bias-circuit impedance of this circuit is shown in Fig. 8(b). It intersects the locus of  $(-)$   $Z_{DDR}$  at 0.42 GHz. Compared with the impedance diagram of Fig. 5, it is seen that this bias circuit is more stable against instabilities than that of Fig. 4(a). However, this bias circuit cannot satisfy the absolutely stable condition in the case of a free-running oscillator, where a waveguide short is connected instead of the cylindrical cavity and the absorber. In this case the circuit  $Q_L$  is 38 and  $f_{max(DDR)}$  is about 3.6 GHz as shown in Fig. 8(b). In order to make the circuit  $Q_L$  higher, a cylindrical cavity has been incorporated and, as a result, the oscillator  $Q$  becomes as high as 400 and  $f_{max(DDR)}$  is decreased to 0.34 GHz as listed in Table I.

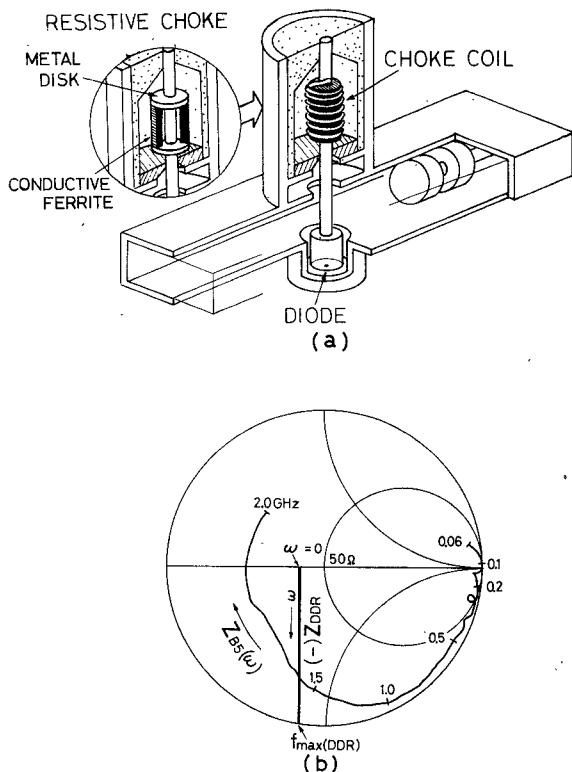


Fig. 10. (a) Modified bias circuit. An RF choke made of a thin wire and a ferrite core is inserted in the center conductor. Insert: Resistive choke which is composed of the lossy conductive ferrite (TDK-X502-RH) and the metal disk. (b) The bias-circuit impedance locus  $Z_{B_s}$  of the circuit using an RF choke. Compared with Fig. 8(b), the impedance has been made fully large at frequencies lower than about 1 GHz.

Thus the  $f_{\max}$  is made lower than the frequency of intersection at 0.42 GHz and the circuit is made stable.

Fig. 9 shows the performance of the developed DDR oscillator with the oscillation frequency of 51.86 GHz. The maximum output of 1 W was achieved at the bias current of 420 mA. The deviation of oscillation frequency was less than 3 MHz within a bias-current range from 220 to 420 mA. This gives a stability factor of 0.3 ppm/mA. Degradation of the oscillation frequency spectrum due to bias-current instabilities was not observed at any bias-current level.

In order to suppress the bias-circuit instability in free-running DDR oscillators without using a high  $Q$  cavity, it is necessary to modify the bias circuit so that the intersecting frequency between the two impedance loci is higher than  $f_{\max(\text{DDR})}$ . Fig. 10(a) shows an example of this scheme. In this circuit, an RF choke made of a thin wire and a ferrite core is inserted in the center conductor. In this configuration, the low-frequency range of instabilities is rejected by the RF choke and the high-frequency range is absorbed in the high impedance coaxial termination. Consequently, the bias-circuit impedance is made high to satisfy the absolutely stable condition. The measured bias-circuit impedance is shown in Fig. 10(b) when the inductance of RF choke was 5  $\mu\text{H}$ . Compared with the impedance diagram of Fig. 8(b), it is seen that a drastic improvement is accomplished in the bias-circuit impedance, and low-frequency instabilities will not occur when  $f_{\max(\text{DDR})}$  is less than 1.5 GHz.

Another method of realizing the resistive choke is also illustrated in the insert in Fig. 10(a). In this method a small tube of lossy conductive ferrite (TDK-X502-RH) is put around the center conductor. The equivalent circuit of this network is a parallel connection of the resistor ( $R \sim 500 \Omega$ ) and the inductor ( $L \sim 0.04 \mu\text{H}$ ). This method has a simpler construction than that using the thin wire and so it is more suitable for millimeter-wave oscillators. It is estimated that free-running DDR oscillators will be operated at their maximum output power without low-frequency instabilities by using these very small stabilizing networks.

## CONCLUSION

The low-frequency instabilities in millimeter-wave DDR IMPATT diodes have been investigated. The equation representing the low-frequency impedance of DDR diodes has been derived. It has been shown that DDR diodes are more prone to instability than SDR diodes. New oscillator circuits with bias-circuit configurations for eliminating the low-frequency instability are proposed. The DDR diodes mounted in these circuits exhibited a maximum oscillation power of 1.6 W at 55.5 GHz with 11.5-percent efficiency. A frequency-stabilized oscillator has also been constructed with maximum output power of 1 W and frequency stability less than 0.3 ppm/mA at 51.86 GHz. It was shown that, by eliminating the low-frequency instability, millimeter-wave DDR IMPATT diodes can practically generate high power with high efficiency.

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